

A Very Large Lepton Collider in a VLHC tunnel

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- Design strategy
- Intensity Limitations
- RF and Optics parameters: Arc, IR
- Lifetime
- Scaling the beam-beam parameter
- Luminosity, Energy reach
- Parameters of a 233 km ring, $\mathcal{L} = 10^{33} \text{cm}^{-2} \text{sec}^{-1}$, $E = 185 \text{ GeV}$
- Low energy design at Z0 mass $E_{cm} = 90 \text{ GeV}$
- Accelerator Physics Challenges

Design Strategy

- Use the maximum RF power available
- Operate at the beam-beam limit

Synchrotron radiation power lost by *both beams*,

$$P_T = 2C_\gamma \frac{E^4 I}{e\rho} \quad (1)$$

$$C_\gamma = (4\pi/3)(r_e/(m_e c^2)^3) = 8.86 \times 10^{-5} \text{ [m/GeV}^3\text{]}.$$

Luminosity

$$\mathcal{L} = \frac{f_{rev} M_b N_b^2}{4\pi \sigma_x^* \sigma_y^*} \quad (2)$$

Vertical beam-beam tune shift

$$\xi_y = \frac{r_e N_b \beta_y^*}{2\pi \gamma \sigma_x^* \sigma_y^*}, \quad \sigma_y^* \ll \sigma_x^* \quad (3)$$

Replacing one power of bunch intensity,

$$\mathcal{L} = \frac{1}{2er_e \beta_y^*} \xi_y \gamma I \quad (4)$$

or

$$\boxed{\mathcal{L} \gamma^3 = \frac{3}{16\pi r_e^2 (m_e c^2)} \frac{\xi_y P_T}{\beta_y^*} \rho} \quad (5)$$

Interpretations

- At fixed \mathcal{L} , P_T and ξ_y

$$E \propto \rho^{1/3}$$

This determines the maximum allowable energy at these parameters.

- At fixed bend radius or circumference C , P_T and ξ_y

$$\mathcal{L} \propto \gamma^{-3}$$

- At constant luminosity, the maximum energy E increases if C , P_T , ξ_y increase and β_y^* decrease.
- ...

Intensity Limitations

Bunch intensity limitations

- At top energy, the limit is set by the beam-beam interactions.

Limits from the desired collisions are included in the design, there may be additional limits from parasitic collisions.

- At injection energy, the transverse mode coupling instability sets the limit. At the threshold the $m = 0$ and $m = -1$ modes of the betatron modes $\omega_\beta + m\omega_s$ become degenerate. Threshold bunch current

$$I_b^{TMCI} \simeq \frac{8f_{rev}\nu_s E}{e \sum_i \beta_i k_{\perp i}(\sigma_s)} \quad (6)$$

$k_{\perp i}$ is a bunch length dependent transverse mode loss factor. At LEP TMCI limits the bunch current to below 1mA. I assume that similar bunch intensities as in LEP will be stable in the large ring but this may be optimistic ...

Beam intensity limitations

- This is primarily determined by the available RF power.
- Cryogenic cooling power.

The dynamic heat load on the cavities includes a contribution from the beam

$$P_{dynamic}^{beam} = 2R_m(\sigma_s)I_b I_e \quad (7)$$

- HOM power in cavities.

Basic Parameters

- Choose β_x^*, β_y^* (limitations determined by IR optics, bunch length)
- Determine the maximum energy E from \mathcal{L} , β_y^* , P_T and ξ_y (choice of ξ_y has to be self-consistent with the energy) for a given circumference.
- Bunch intensity N_b is set by TMCI limitations
- Choose a coupling ratio (determined by β_y^*/β_x^*)

$$\kappa = \frac{\epsilon_y}{\epsilon_x} \quad (8)$$

- Equilibrium emittance is found from

$$\epsilon_x = \frac{N_b}{\gamma \xi_y} \left(\frac{r_e}{2\pi} \sqrt{\frac{\beta_y^*}{\kappa \beta_x^*}} \right) \quad (9)$$

where factors within () are assumed to stay constant.

- Choose a phase advance per cell μ_C (upper limit usually determined by chromaticity sextupoles).
- The cell length L_C is determined by the equilibrium emittance

$$\epsilon_x \approx \left(\frac{C_q R}{J_x \rho} \left[\frac{L_C}{\mu_c} \right]^3 \right) \frac{\gamma^2}{R^3} \quad (10)$$

$$C_q = 55\hbar c / (32\sqrt{3}(m_e c^2)) = 3.83 \times 10^{-13} \text{ [m]}$$

- Filling factors f_1 and f_2

$$R = f_1 \frac{C}{2\pi}, \quad \text{and} \quad \rho = f_2 R, \quad f_1, f_2 < 1 \quad (11)$$

R is the arc radius, ρ is the bend radius. Typically $1 < R/\rho \leq 1.25$.

- Maximum number of bunches is determined by the beam power P_T

$$M_b^{max} = \left(\frac{P_T}{2C\gamma} \right) \frac{\rho}{f_{rev} N_b E^4} \quad (12)$$

RF parameters

Requirements

- The RF must replenish the energy lost per turn.
- The RF must provide an acceptable quantum lifetime.

Energy Gain

$$eV_{RF} \sin \phi_s = U = C_\gamma \frac{E^4}{\rho} \quad (13)$$

The longitudinal quantum lifetime is determined by the energy headroom $N_{QL} = \Delta E_{RF} / \sigma_E$ as

$$\tau_{quant;s} = \frac{\tau_s}{N_{QL}^2} \exp\left[\frac{1}{2} N_{QL}^2\right] \quad (14)$$

$$\sqrt{\frac{1}{\pi h \eta_{slip}} e V_{RF} E G(\phi_s)} = N_{QL} \sqrt{\frac{C_q}{J_s \rho} \frac{E^2}{m_e c^2}} \quad (15)$$

where

$$G(\phi_s) = 2 \cos \phi_s - (\pi - 2\phi_s) \sin \phi_s \quad (16)$$

Typically $N_{QL} \sim 10$. The two requirements determine the equation for the synchronous phase ϕ_s and the RF voltage V_{RF} .

RF frequency

- The RF acceptance $(\Delta E/E)_{accep} \propto 1/\sqrt{h}$ so lower RF frequencies increase the acceptance.
- However high power klystrons are cheaper above frequencies of 300MHz.

LEP operates with 352MHz. For this design we chose an RF frequency of 400MHz.

Arc parameters: phase advance and cell length

Equilibrium emittance

- The emittance *decreases* as the phase advance *increases*, reaching a minimum at 135° . In a lattice with FODO cells,

$$\epsilon_x(\mu_x^C) = 4 \frac{C_q \gamma^2}{J_x} \theta^3 \frac{1 - \frac{3}{4} \sin^2(\mu_x^C/2) + \frac{1}{60} \sin^4(\mu_x^C/2)}{\sin^2(\mu_x^C/2) \sin \mu_x^C}. \quad (17)$$

but

- Stronger focusing increases the chromaticity and the strength of the chromaticity sextupoles which can limit the dynamic aperture.

Typically

$$60^\circ \leq \mu_c < 120^\circ$$

For example, LEP has operated with $(60^\circ, 60^\circ)$ at 45GeV, and since then $(90^\circ, 60^\circ)$, $(90^\circ, 90^\circ)$ and $(102^\circ, 90^\circ)$ at higher energies.

TMCI threshold

$$I_{thresh}^{TMCI} \propto \frac{\nu_s}{\langle \beta \rangle} \propto \frac{1}{L_c} \cos\left(\frac{\mu_c}{2}\right) \quad (18)$$

The TMCI threshold *increases* if the cell length L_C and phase advance per cell μ_C *decrease*.

Emittance control by changing the RF frequency.

$$\frac{dJ_x}{d\delta} = -\frac{dJ_s}{d\delta} = -4 \frac{L_D}{L_Q} \left[\frac{2 + \frac{1}{2} \sin^2 \mu_C/2}{\sin^2 \mu_C/2} \right] \quad (19)$$

L_D : length of dipoles in a half cell, L_Q : length of a quadrupole.

Required RF frequency shift is related to the momentum deviation δ by

$$\frac{\Delta f_{RF}}{f_{RF}} = -\frac{\Delta R}{R} = -\alpha_C \delta \quad (20)$$

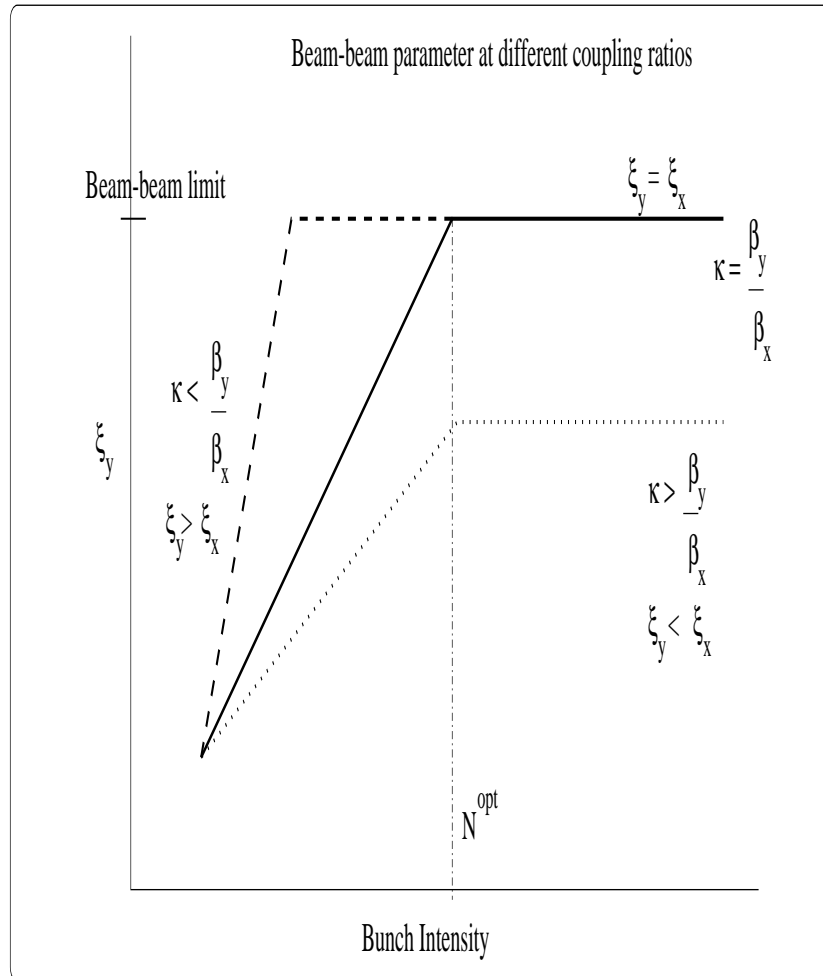
Important to keep ΔR small to minimize loss in physical aperture and transverse quantum lifetime, i.e. design $\Delta J_x / \Delta R$ to be large. This requires lower μ_C and L_Q / L_D to be small i.e. *weaker focusing*.

Example: $C = 228\text{km}$, $L_D = 103.7\text{m}$, $L_Q = 4.1\text{m}$, $\mu_C = 90^\circ$, $\alpha_C = 0.28 \times 10^{-4}$,

$$\frac{\Delta J_x}{\Delta R} = 0.54 \text{ / [mm]}$$

This is large enough to be useful.

Beam-beam limit



IR Parameters

Lower limits on β^*

- Set by the tolerable beam size in the IR quadrupoles and the chromaticity of these quadrupoles. With $\beta_y^* \ll \beta_x^*$, aperture and chromaticity limitations will first arise in the vertical plane.
- $\beta_y^* \gg \sigma_s$ to prevent the luminosity loss due to the hourglass effect.

β^* , coupling and the beam-beam limit

Beyond the beam-beam limit, $\epsilon_x \propto I$, $\xi_x, \xi_y \sim \text{const.}$ and $\mathcal{L} \propto I$.

$$\xi_x = \left[\sqrt{\frac{\kappa}{\beta_y^*/\beta_x^*}} \right] \xi_y, \quad \mathcal{L} \propto \xi_y \quad (21)$$

If $\kappa > \beta_y^*/\beta_x^*$, then $\xi_x > \xi_y$, the beam-beam limit is reached first in the horizontal plane. ξ_y never reaches its maximum value and since $\mathcal{L} \propto \xi_y$ the maximum luminosity is not obtained. So require $\kappa \leq \beta_y^*/\beta_x^*$ or $\xi_y \geq \xi_x$.

Optimal coupling: $\kappa = \beta_y^*/\beta_x^*$ and $\xi_x = \xi_y$.

$$N^{opt} = \frac{2\pi\gamma\epsilon_x}{r_e} \xi_y \quad (22)$$

If $\kappa < \beta_y^*/\beta_x^*$, the limit is reached at intensity

$$N^{limit} = \frac{2\pi\gamma\epsilon_x}{r_e} \sqrt{\frac{\kappa}{\beta_y^*/\beta_x^*}} \xi_y < N^{opt} \quad (23)$$

Beam Lifetime

- Radiative Bhabha scattering process $e^+e^- \rightarrow e^+e^-\gamma$.

$$\begin{aligned}
 \tau_L &= \frac{1}{N_{IP}} \frac{M_b N_b}{\mathcal{L} \sigma_{e^+e^-}} \\
 &= \left[\frac{2r_e \beta_y^*}{N_{IP} \xi_y \sigma_{e^+e^-}} \right] \frac{1}{\gamma f_{rev}} \\
 &\propto \frac{1}{\gamma \xi_y}
 \end{aligned} \tag{24}$$

The cross-section $\sigma_{e^+e^-}$ depends on the energy acceptance and has a weak logarithmic dependence on energy. In the energy range from 175 - 200 GeV per beam, $\sigma_{e^+e^-} \sim 0.36$ mbarns assuming an RF acceptance of 1%.

- Beam-gas scattering.
- Compton scattering off thermal photons.

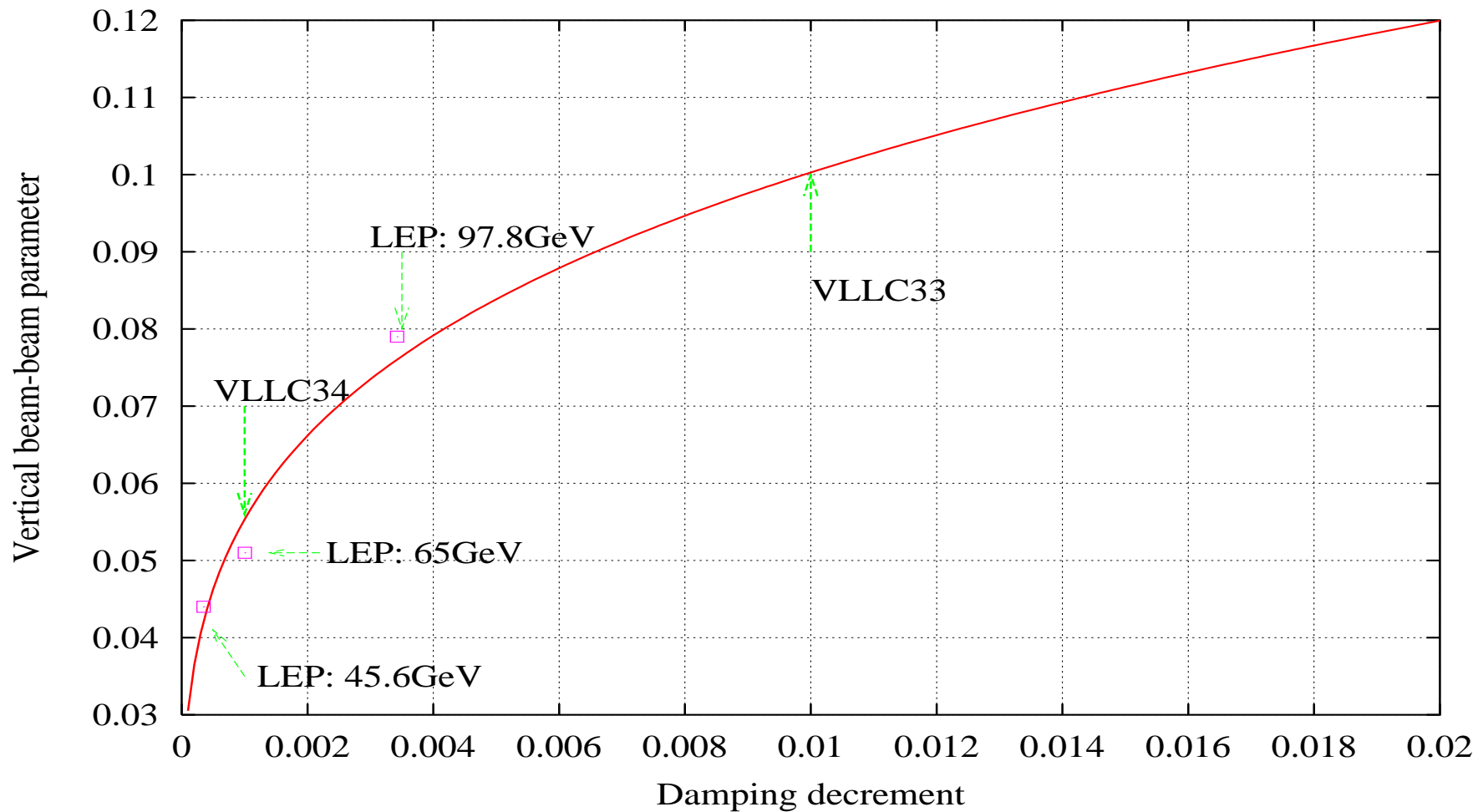
Total lifetime

$$\frac{1}{\tau} = \sum_i \frac{1}{\tau_i} \tag{25}$$

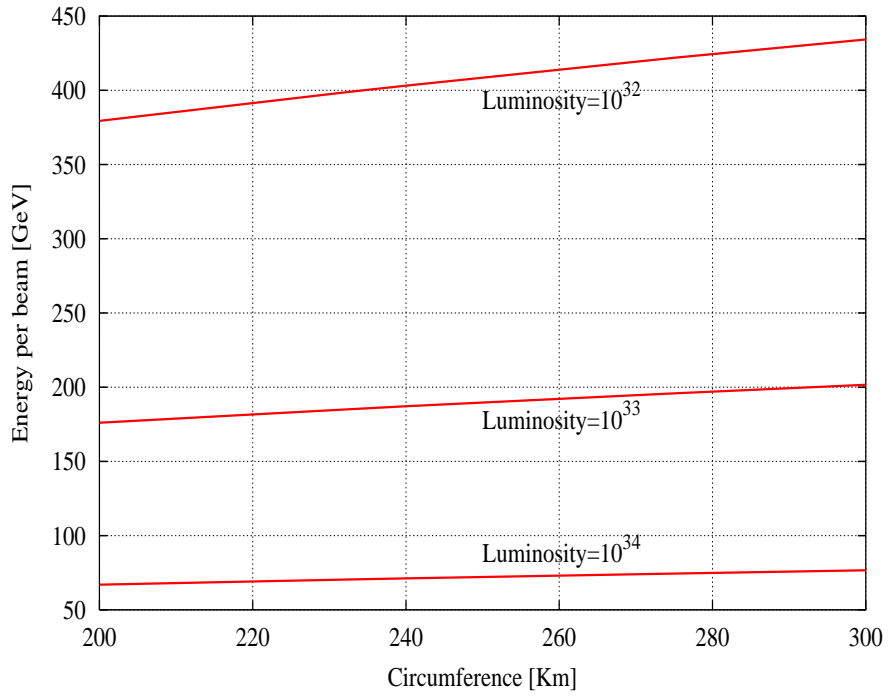
Example: LEP

| Process | τ [hrs] |
|--|--------------|
| Radiative Bhabha scattering | 5.8 |
| Compton scattering | 60 |
| Beam-gas scattering (pressure=0.6 nTorr) | 80 |
| Total | 5.0 |

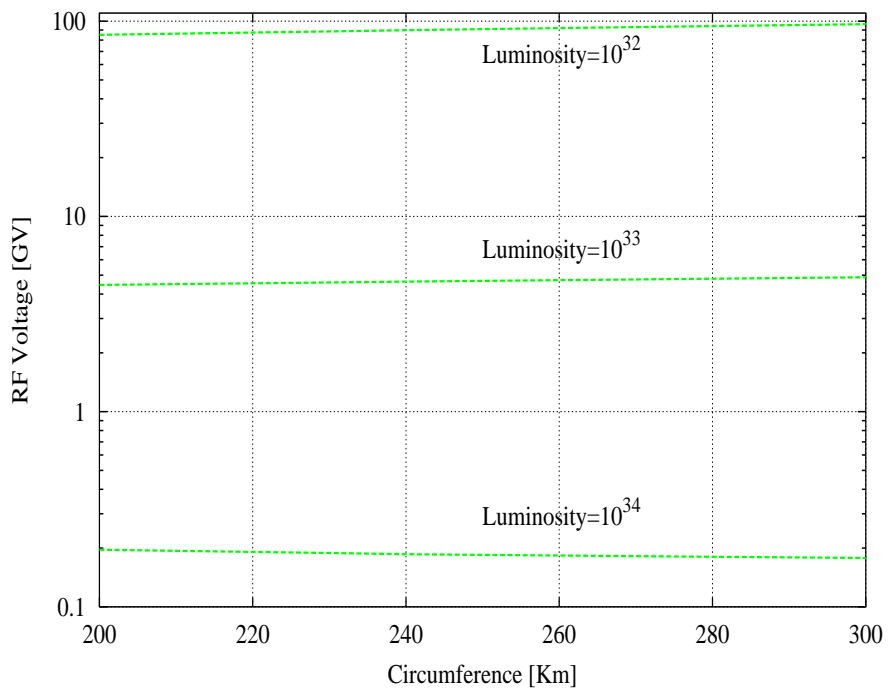
Scaling of the beam-beam parameter



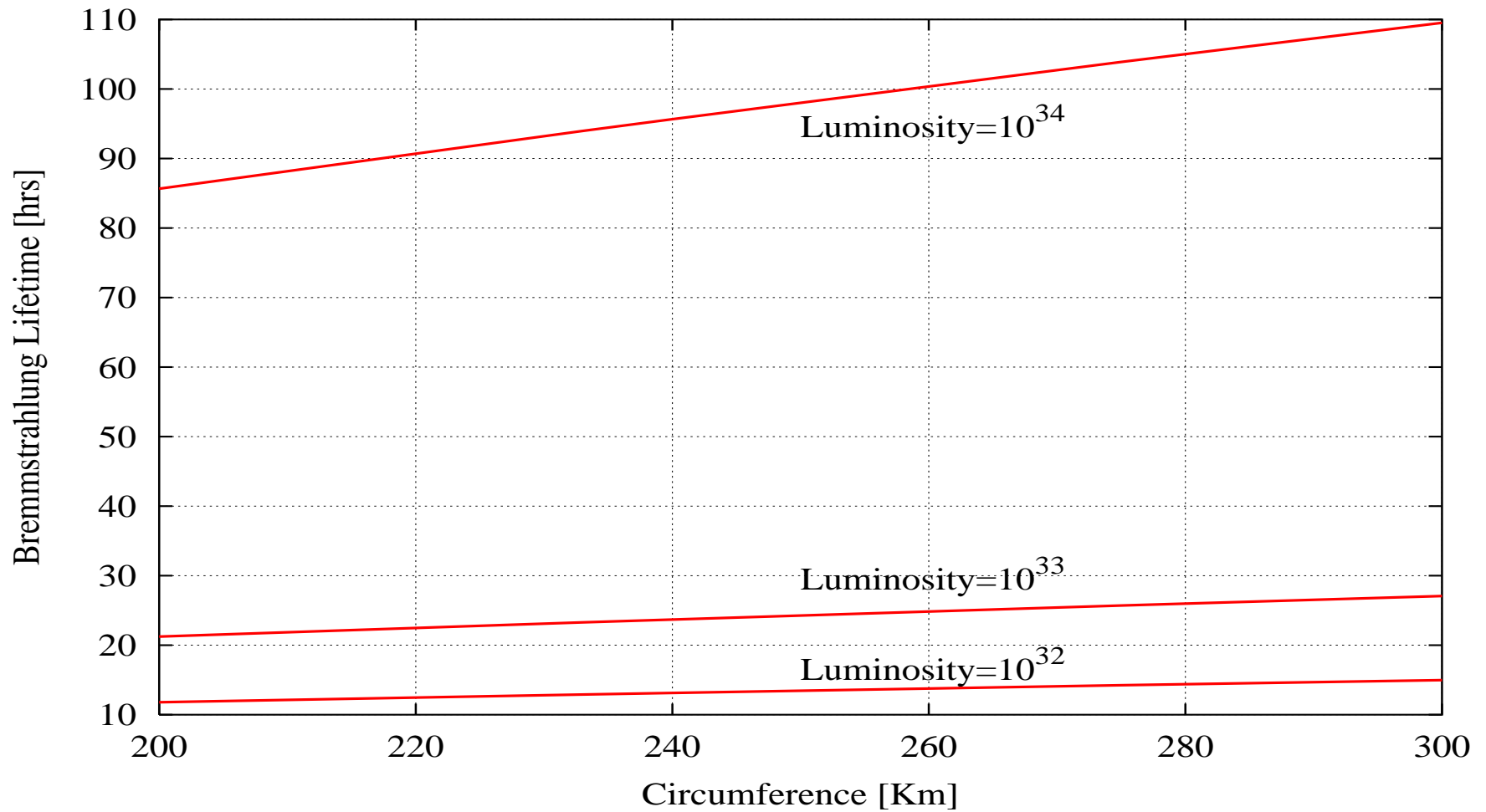
Energy vs Circumference: synch. rad. power = 100MW



RF Voltage vs Circumference: synch. rad. power = 100MW



Luminosity lifetime vs Circumference: synch. rad. power = 100MW



$e^+ - e^-$ Collider Parameters

| | |
|---|------------------------|
| Energy [GeV] | 185.00 |
| Luminosity | 1.00×10^{33} |
| Synch. radiation power(both beams) [MW] | 100.712 |
| σ_x^*, σ_y^* [microns] | 77.520, 3.876 |
| Number of bunches | 126 |
| Bunch spacing [km] | 1.8492 |
| Particles per bunch | 4.851×10^{11} |
| Bunch current [mA] | 0.100 |
| Emittances [nano-m] | 6.009, 0.300 |
| Beam-beam parameter | 0.10650 |
| Damping decrement | 0.01080 |
| Single beam current [mA] | 12.600 |
| Brho [Tesla-m] | 617.078 |
| Arc tune | 215.356 |
| Phase advance per cell [deg] | 90.000 |
| Dipole field [T] | 0.02376 |
| Focal length of cell [m] | 80.025 |
| Quad gradient [T/m] | 15.595 |
| Quad field at $1\sigma_x^{max}$ /dipole field | 1.000 |
| Cell: β^{max}, β^{min} [m] | 386.395, 66.295 |
| Cell: $\sigma_x^{max}, \sigma_x^{min}$ [mm] | 1.524, 0.631 |
| Cell: $\sigma_y^{max}, \sigma_y^{min}$ [mm] | 0.341, 0.141 |
| Max apertures required [cm] | 2.524, 1.341 |
| Max and min disp. [m] | 1.117, 0.534 |
| Momentum compaction | 0.2226E-04 |
| Energy loss per turn [GeV] | 3.996 |
| Damping time [turns] | 46 |
| RF Voltage [GV] | 4.57254 |
| Synchronous phase [deg] | 60.930 |
| Relative energy spread | 0.9834E-03 |
| RF acceptance | 0.9822E-02 |
| Synchrotron tune | 0.11501 |
| Bunch length [mm] | 7.058 |
| Longitudinal emittance [eV-sec] | 0.01346 |
| Bremm cross-section [barns] | 0.36022 |
| Bremm lifetime [hrs] | 23.6 |
| Polarization time [hrs] | 2.215 |
| Critical energy [keV] | 452.611 |
| Critical wavelength [A] | 0.023 |
| Number of photons/m/sec | 0.115E+17 |
| Gas load [torr-L/m-sec] | 0.104E-07 |
| Linear Power load(both beams) [kW/m] | 0.517 |
| specif press. rise [Torr/mA] | 0.138E-12 |
| specific current [mA] | 7238.604 |
| spec. current/beam current | 574.49 |

Comments on Parameters

- $C(VLLC33)/C(LEP) \sim 8.5$, $P_T(VLLC)/P_T(LEP) \sim 7$
 $\Rightarrow \mathcal{L}(VLLC33)/\mathcal{L}(LEP) \sim 10$ at almost double the energy.
- The $e^+ - e^-$ bremsstrahlung lifetime in VLLC33 is significantly longer at 23 hours.
- The beam sizes in the two machines are comparable. Hence vacuum chamber dimensions in VLLC33 can be similar to those in LEP if a two-ring machine (long-range interactions in a single ring machine may require a larger aperture).
- $B_{dip}(VLLC)/B_{dip}(LEP) \sim 1/5$. Warm iron magnets will suffice. But good shielding from stray magnetic fields will be more important.
- The quadrupole gradient is determined by requiring that synchrotron radiation in quadrupoles be small. Require $B_{quad}(r = 1\sigma)/B_{dip} = 1$ (E. Keil).
- The critical energy is smaller in VLLC33 so shielding against synchrotron radiation as in LEP should be adequate for VLLC33. The photon flux per unit length is almost the same in the two machines.
- The RF voltage required for VLLC33 is higher at 4.7GV compared to 3.1GV for LEP.
- $f_1 = f_2 = 0.84$ was chosen to have the same ratio $\rho : C/(2\pi)$ as in LEP. A more aggressive choice of $2\pi\rho/C = 0.81$ yields e.g. maximum energy $E_{max} = 193\text{GeV}$, RF voltage $V_{RF} = 4883\text{MV}$.
- We chose optimum coupling, i.e. $\epsilon_y/\epsilon_x = \beta_y^*/\beta_x^*0.05$ which $\Rightarrow \xi_x = \xi_y$. If we reduce it to $\epsilon_y/\epsilon_x = 0.025$, then $\xi_x = 0.071$, $\xi_y = 0.1$. Optics and beam size parameters change, e.g. $\epsilon_x = 11.8\text{nm}$, cell length=278m, $\beta^{max} = 475\text{m}$, $D_x^{max} = 1.72\text{m}$, $\sigma_x^{max} = 2.4\text{mm}$, $\nu_s = 0.156$, $\sigma_l = 8.1\text{mm}$. The RF voltage increases to 4780MV, most other parameters are relatively unaffected.
- We chose $N_{QL} = 10$ to ensure sufficient quantum lifetime. At LEP $N_{QL} = 6.6$. If we assume this value for the 228km ring, the RF voltage is lowered from 4.66GV to 4.43GV.

Comparison of parameters

| Parameter | Bunch current $I_b = 0.1\text{mA}$ | | | Bunch current $I_b = 0.05\text{mA}$ | |
|---|------------------------------------|-----------------------------------|--------------------------------|-------------------------------------|--------------------------------|
| | $\kappa = \beta_y^*/\beta_x^*$ | $\kappa = 0.5\beta_y^*/\beta_x^*$ | $\kappa = \beta_y^*/\beta_x^*$ | $\kappa = \beta_y^*/\beta_x^*$ | $\kappa = \beta_y^*/\beta_x^*$ |
| | $\xi_y = 0.1$ | $\xi_y = 0.1$ | $\xi_y = 0.08$ | $\xi_y = 0.1$ | $\xi_y = 0.08$ |
| Energy [GeV] | 185 | 185 | 172 | 185 | 172 |
| Emittances ϵ_x, ϵ_y [nm] | 6.1, 0.3 | 8.6, 0.21 | 8.1, 0.40 | 3.0, 0.15 | 4.1, 0.2 |
| Number of bunches | 123 | 123 | 167 | 246 | 332 |
| τ_{bremm} [hrs] | 23 | 23 | 31 | 23 | 31 |
| Arc tune | 215 | 191 | 186 | 271 | 234 |
| Cell Length [m] | 227 | 255 | 262 | 180 | 209 |
| Arc $\sigma_x^{max}, \sigma_x^{min}$ [mm] | 1.53, 0.63 | 1.93, 0.80 | 1.90, 0.79 | 0.96, 0.40 | 1.2, 0.5 |
| Synchrotron tune ν_s | 0.116 | 0.135 | 0.128 | 0.084 | 0.094 |

The luminosity is $10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$ in each case. Ring circumference is 233 km and the synchrotron radiation power is fixed at 100 MW.

Low Energy Operation

At “low energies”, the ring is

- not limited by available power
- only constrained by the beam-beam tune shift

$$\begin{aligned}\mathcal{L} &= \frac{\pi}{r_e^2} M_B f_{rev} \left[\frac{\sigma_x^* \sigma_y^*}{(\beta_y^*)^2} \right] \gamma^2 \xi_y^2 \\ &= \frac{\pi}{r_e^2} M_B f_{rev} \left[\frac{\kappa \beta_x^*}{(\beta_y^*)^3} \right]^{1/2} \gamma^2 \xi_y^2 \epsilon_x\end{aligned}$$

In this regime,

Luminosity increases with the emittance $\mathcal{L} \propto \epsilon_x$

This requires *filling the aperture* at these energies.

Sufficient physical aperture

$$\text{Aperture} \approx 10 * [\sigma_x^2 + (D_x \delta_p^2)]^{1/2} + 1\text{cm(c.o.d)} \quad (26)$$

The bunch intensity is low in order to limit the beam-beam tune shift. Number of bunches have to be increased to increase the luminosity.

What is the minimum bunch spacing at these intensities?

Assume (without justification) a minimum spacing of 5m.

Low Energy Design Strategy

- Increase the emittance by lowering the phase advance per cell μ_C .

- The emittance

$$\epsilon_x = \left(\frac{C_q R}{J_x \rho}\right) \left[\frac{L_C}{R \mu_C}\right]^3 \gamma^2$$

- Find the smallest phase advance so that

$$10 * [\sigma_x^2 + (D_x^{max} \delta_p)^2]^{1/2} + 1\text{cm} \leq \text{Aperture}$$

- Find the bunch intensity from the beam-beam tune shift

$$N_b = \left(\frac{2\pi}{r_e} \sqrt{\frac{\kappa}{\beta_y^*/\beta_x^*}}\right) \gamma \epsilon_x \xi_y$$

Check that $N_b \leq N_b^{TMCI}$

- Find the number of bunches M_B from the minimum bunch spacing

$$M_B f_{rev} = \frac{c}{S_B}$$

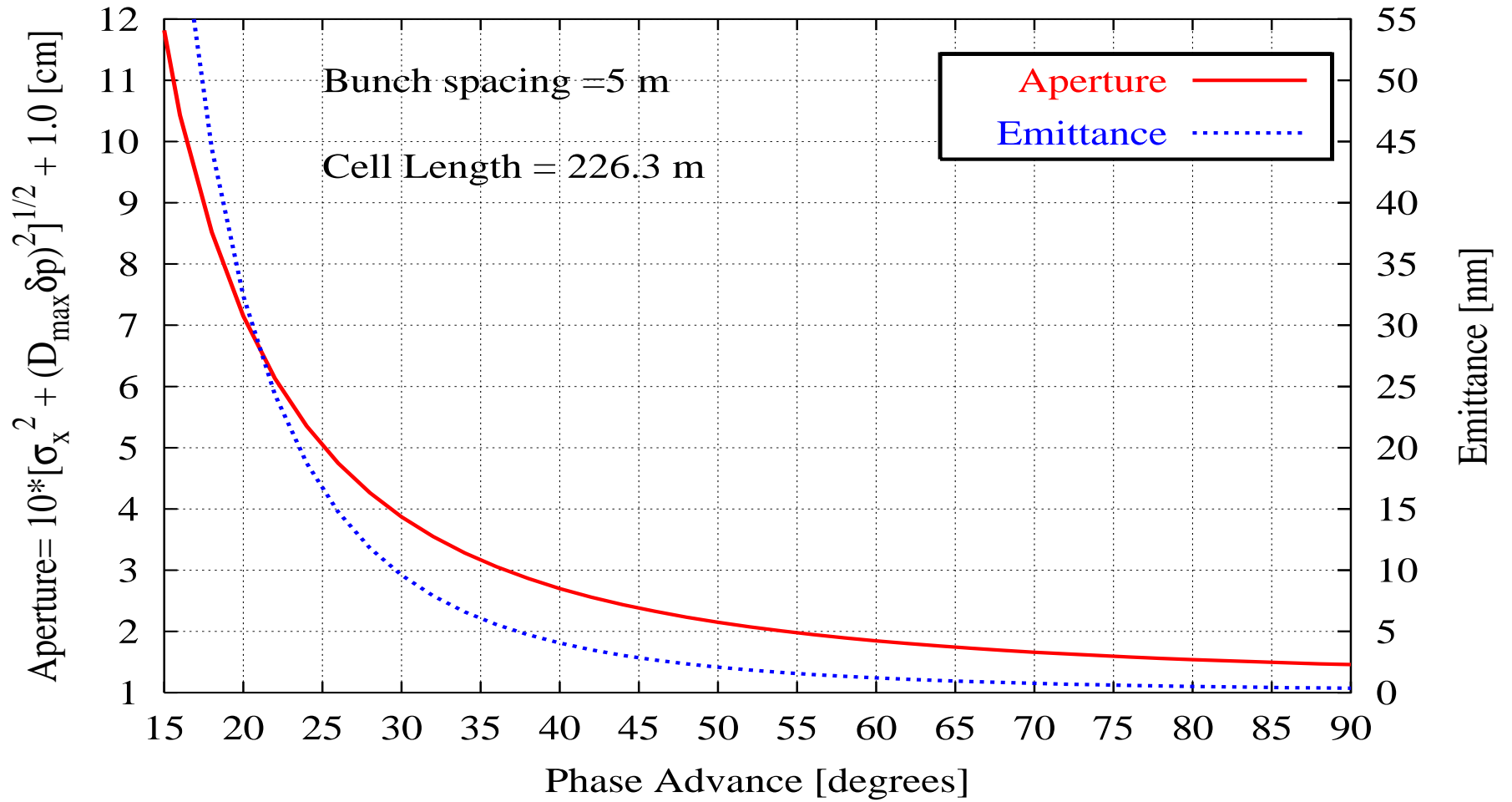
- Luminosity

$$\mathcal{L} = \frac{\pi}{r_e^2} M_B f_{rev} \left[\frac{\kappa \beta_x^*}{(\beta_y^*)^3}\right]^{1/2} \gamma^2 \xi_y^2 \epsilon_x$$

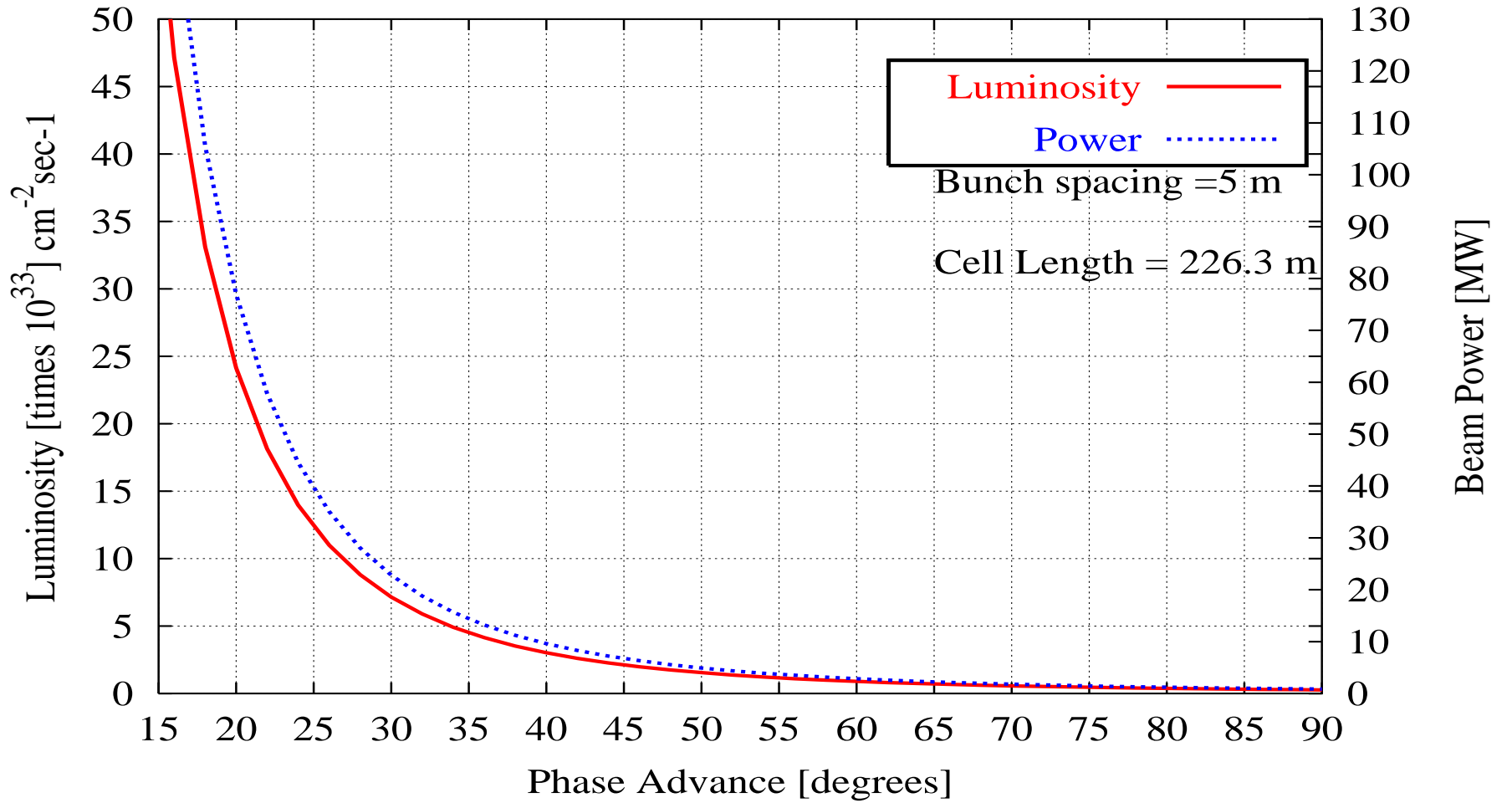
- Alternative strategy of increasing the emittance:

Double the cell length by turning off half the quadrupoles.

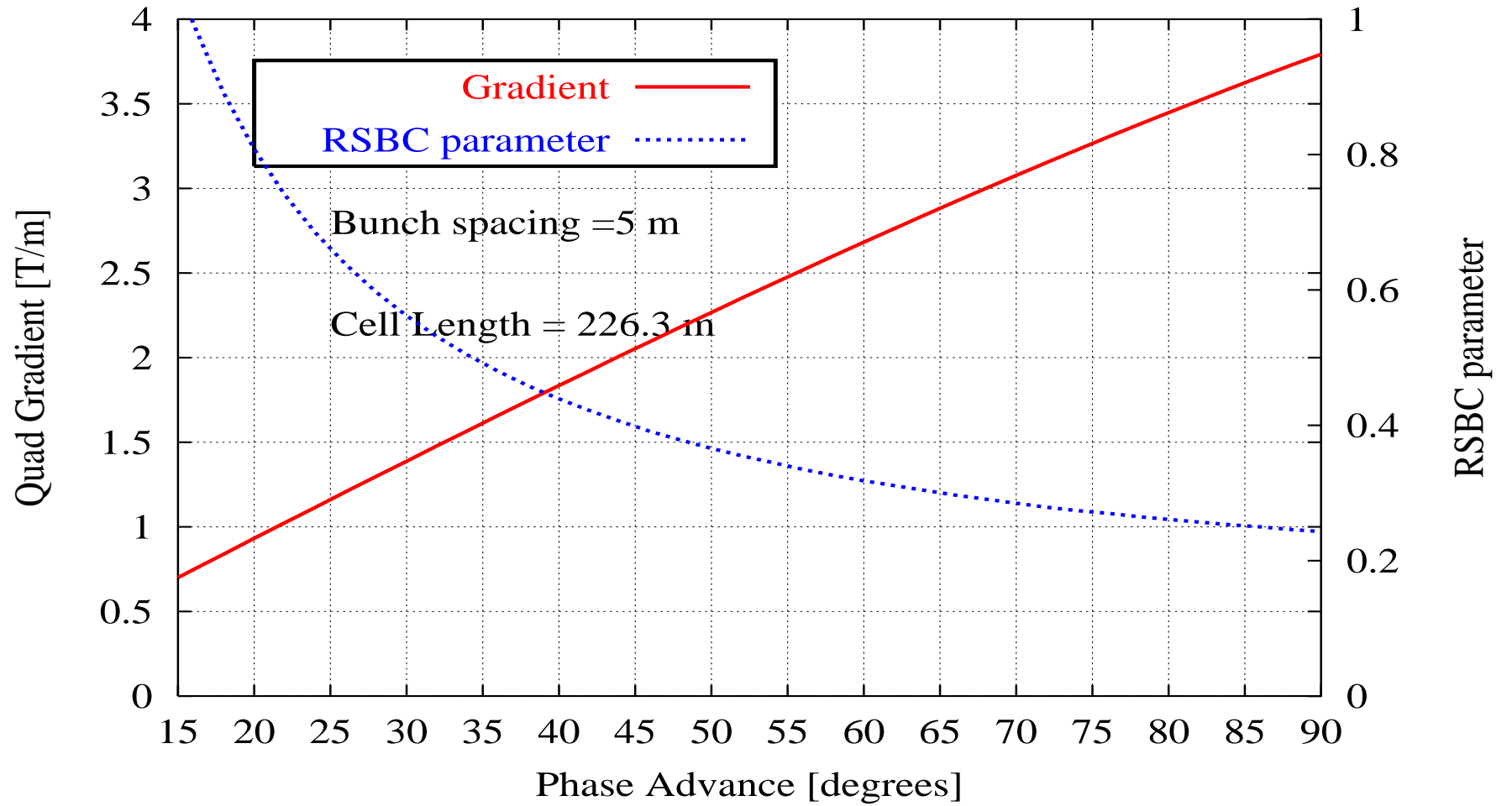
Circumference=233km, Energy = 45 GeV



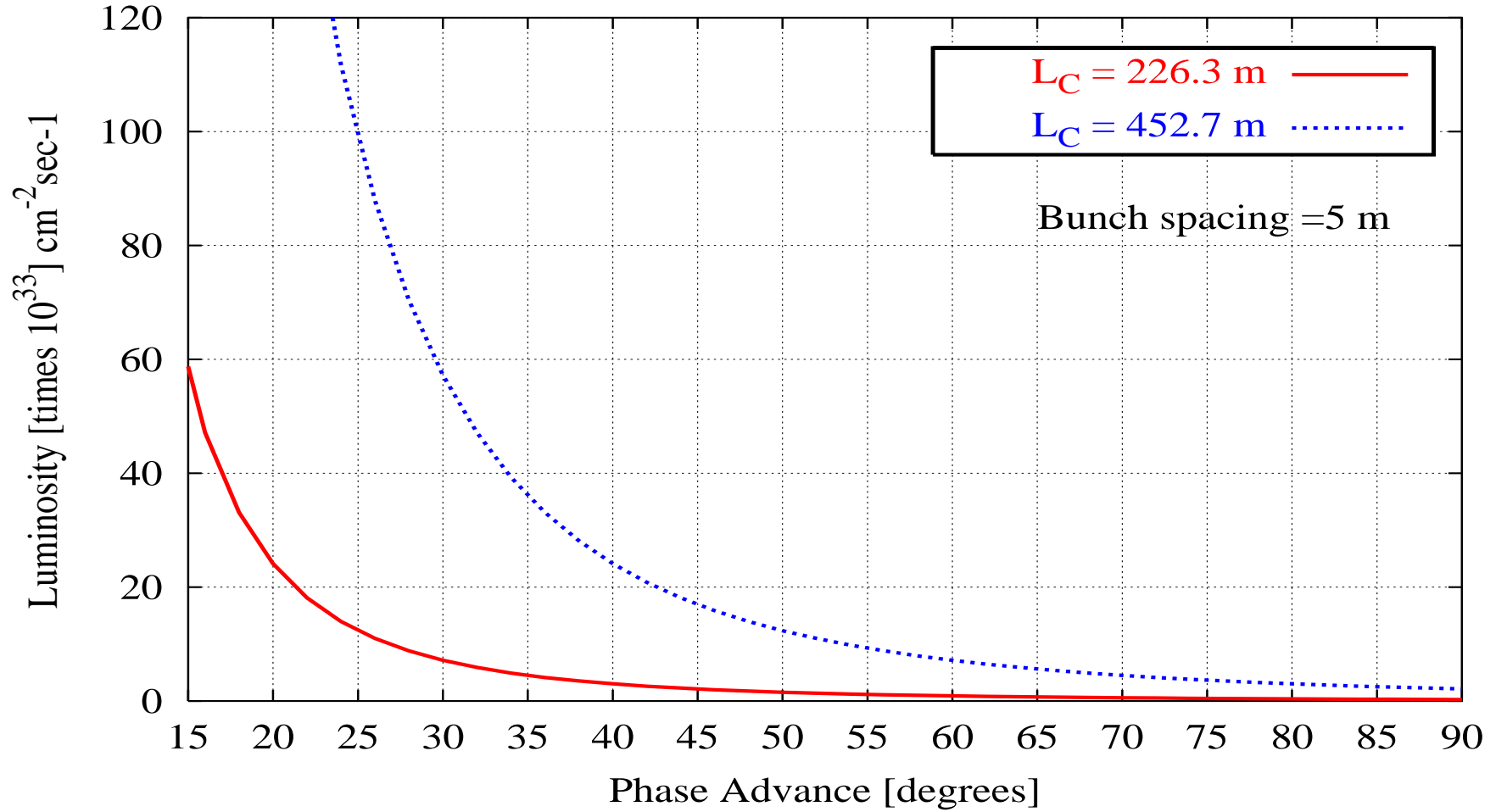
Circumference=233km, Energy = 45 GeV



Circumference=233km, Energy = 45 GeV

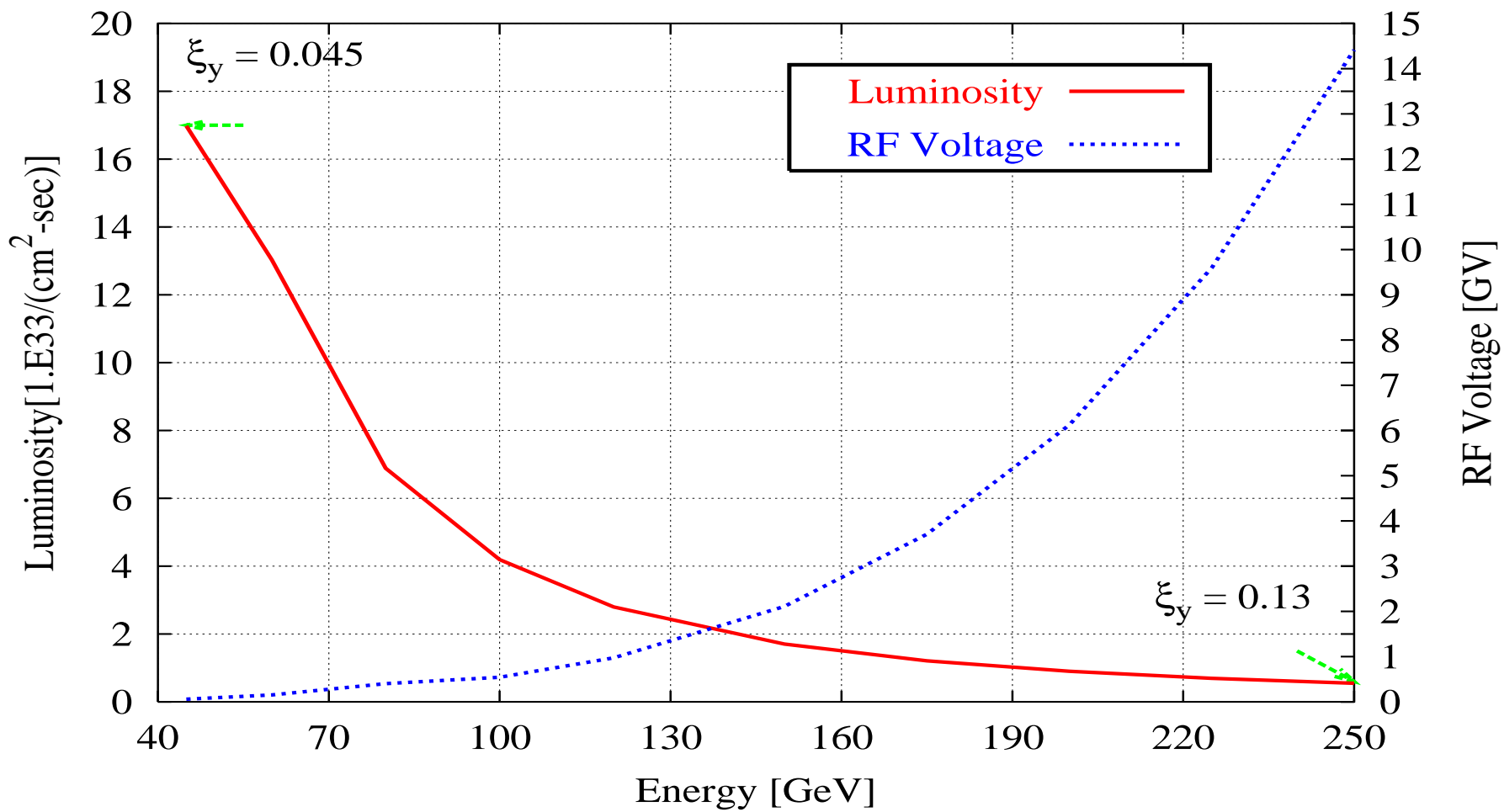


Circumference=233km, Energy = 45 GeV



| | |
|---|------------------------|
| Energy [GeV] | 45.00 |
| Luminosity | 16.98×10^{33} |
| Synch. radiation power(both beams) [MW] | 54.046 |
| σ_x^*, σ_y^* [microns] | 150.849, 7.542 |
| Number of bunches | 46599 |
| Bunch spacing [km] | 0.0050 |
| Particles per bunch | 2.011×10^{11} |
| Bunch current [mA] | 0.04145 |
| Emittances [nano-m] | 22.755, 1.138 |
| Beam-beam parameter | 0.04500 |
| Damping decrement | 0.00016 |
| Single beam current [mA] | 1931.485 |
| Brho [Tesla-m] | 150.100 |
| Arc tune | 53.839 |
| Phase advance per cell [deg] | 22.500 |
| Dipole field [T] | 0.00578 |
| Focal length of cell [m] | 290.051 |
| Quad gradient [T/m] | 1.047 |
| Quad field at $1\sigma_x^{max}$ /dipole field | 0.726 |
| Cell: $\beta_x^{max}, \beta_x^{min}$ [m] | 706.857, 476.078 |
| Cell: $\sigma_x^{max}, \sigma_x^{min}$ [mm] | 4.011, 3.291 |
| Cell: $\sigma_y^{max}, \sigma_y^{min}$ [mm] | 0.897, 0.736 |
| Max apertures required [cm] | 5.011, 1.897 |
| Max and min disp. [m] | 11.902, 9.786 |
| Momentum compaction | 0.2924E-03 |
| Energy loss per turn [GeV] | 0.0140 |
| Damping time [turns] | 3216 |
| RF Voltage [GV] | 0.05729 |
| Synchronous phase [deg] | 14.136 |
| Relative energy spread | 0.2392E-03 |
| RF acceptance | 0.2400E-02 |
| Synchrotron tune | 0.13365 |
| Bunch length [mm] | 19.409 |
| Longitudinal emittance [eV-sec] | 0.00219 |
| Bremm cross-section [barns] | 0.45368 |
| Bremm lifetime [hrs] | 168.9 |
| Polarization time [hrs] | 2600.792 |
| Critical energy [keV] | 6.514 |
| Critical wavelength [A] | 1.593 |
| Number of photons/m/sec | 0.430E+18 |
| Gas load [torr-L/m-sec] | 0.387E-06 |
| Linear Power load(both beams) [kW/m] | 0.277 |
| specif press. rise [Torr/mA] | 0.336E-13 |
| specific current [mA] | 29758.71 |
| spec. current/beam current | 15.41 |

Circumference=233km, synch. rad. power = 100MW



Stored Energy, Site Power

Stored energy per beam

- LEP (E=98GeV, I=2.88mA, $f_{rev}=11.25\text{kHz}$)
Stored energy $\sim 25\text{kJ}$
- VLLC (E=185GeV, I=12.4mA, $f_{rev}=1.315\text{kHz}$)
Stored energy $\sim 1.74\text{MJ}$
- VLHC (E=7TeV, I=0.53A, $f_{rev}=11.25\text{kHz}$)
Stored energy $\sim 334\text{MJ}$

Stored energy is not large in comparison to hadron colliders.

Site Power requirements

- Beam power 100MW. Assuming a klystron efficiency of 50%, this requires a wall plug power of 200MW.
- Cryogenic cooling power $\sim 100\text{ MW}$ (?).
- 100 MW of heat has to be extracted by cooling water. Some fraction of this power will be required to pump the water.

Power requirements are significant.

Accelerator Physics Challenges

- Combating TMCI

The transverse impedance of the beam pipe alone is close to the threshold impedance. Other major contributions to the impedance from cavities, bellows, ... will likely increase the impedance to beyond threshold.

Possible solutions

- Coalescing bunches at top energy
- Feedback system
- Raising the injection energy
- Increasing the bunch length and the synchrotron tune.
- ...

- If a single ring machine, then long-range beam-beam interactions with many bunches will limit the beam stability. Will likely affect the required physical aperture.
- With many bunches, multi-bunch instabilities may be an issue.
- Avoiding synchro-betatron resonances e.g. those driven by dispersion in the cavities.
- The beam-beam limit may not increase with damping decrement as hoped for. In that case, achievable luminosities will be lower.
- At low energy (45 GeV), the beam current is high (~ 2 A) and the number of bunches is very large (about 47000). Operation in a single ring machine at these parameters will be extremely difficult if not impossible and will be challenging even in a two ring machine of this circumference.
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