

Hyperon Decay 2007

- CP violation in kaon decay measured
 - both ϵ and ϵ' have been measured
- Rare kaon decay $K \rightarrow \pi \nu \nu$ program ongoing
- Hyperon weak decay occurs through the same effective transition as kaon weak decay:

$$s \rightarrow d$$

- So why is hyperon decay still interesting?

- Because the hyperon decay amplitude is simultaneously sensitive to both parity conserving and parity violating weak interactions.
- CP violation in hyperon decay within the SM is correlated to ϵ' and it is small.
- CP violation in hyperon decay beyond SM is limited by theoretical uncertainty in ϵ , and can be much larger.

CP Violation in hyperon decay

- Compare hyperon and anti-hyperon decay
- Each decay of the form $B' \rightarrow B \pi$ has s-wave and p-wave amplitudes which could be $\Delta I=1/2$ or $\Delta I=3/2$.
- The angular distribution with polarized B' is

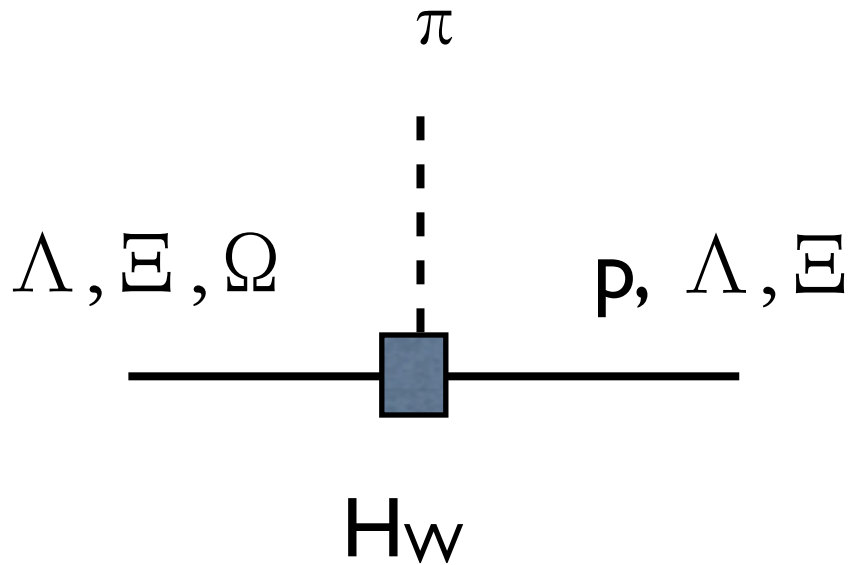
$$\Lambda \rightarrow p\pi^-: \quad \frac{dN(p)}{d\cos\theta} = \frac{N_0}{2}(1 + \alpha_\Lambda P_\Lambda \cos\theta) \quad \alpha = \frac{2\text{Re}(S^*P)}{|S|^2 + |P|^2}$$

- Two CP odd observables are

$$\Delta \equiv \frac{\Gamma - \bar{\Gamma}}{\Gamma + \bar{\Gamma}} \sim \sqrt{2} \frac{S_3}{S_1} \sin(\delta_3^S - \delta_1^S) \sin(\phi_3^S - \phi_1^S)$$

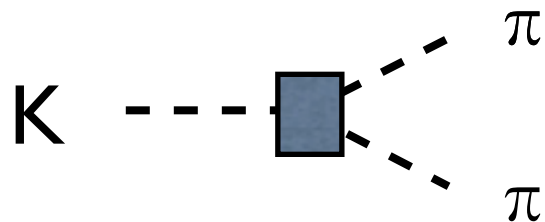
$$A(\Lambda_-^0) \equiv \frac{\alpha + \bar{\alpha}}{\alpha - \bar{\alpha}} \sim -\sin(\delta_1^P - \delta_1^S) \sin(\phi_1^P - \phi_1^S) \sim 0.12 \sin(\phi_1^P - \phi_1^S)$$

S (or D) waves

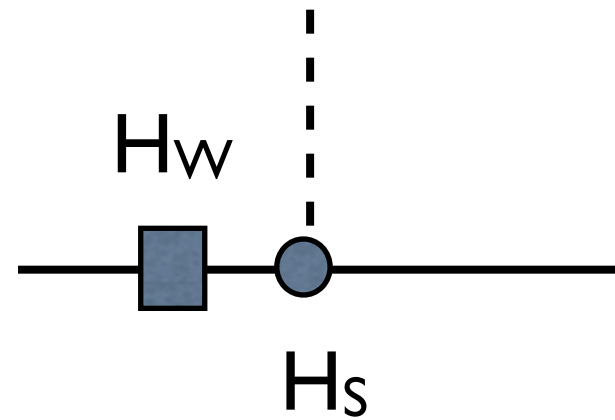


H_W parity violating

CP constraint from ϵ'



P waves

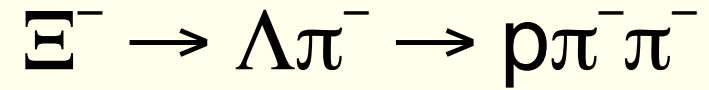


parity conserving

constraint from ϵ



HyperCP looked for asymmetry in

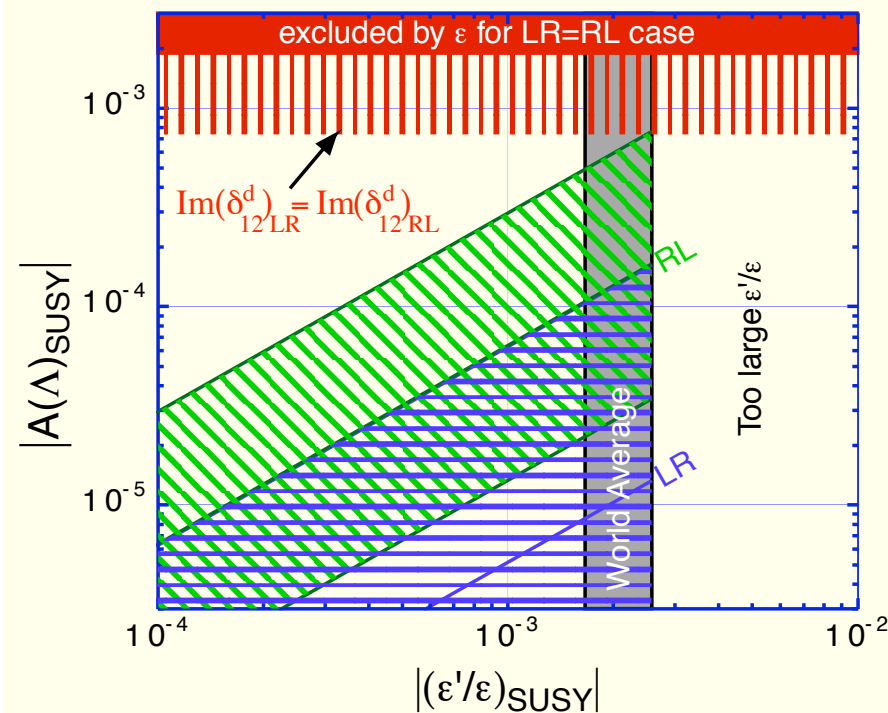


Most recent SM prediction

$$-0.5 \times 10^{-4} < A_{\Xi\Lambda} < +0.5 \times 10^{-4}$$

Beyond the SM, bounds from ϵ allow up to 10^{-3}

SUSY example:



He et al., PRD 61 (2000) 071701(R).

HyperCP result

Phys.Rev.Lett.93:262001,2004.

$$A_{\Xi\Lambda} = [0.0 \pm 5.1(\text{stat}) \pm 4.4(\text{syst})] \times 10^{-4}$$

CP Violation in Ω Decay

$$i\mathcal{M}_{\Omega^- \rightarrow \Xi \pi} = G_F m_\pi^2 \bar{u}_\Xi \mathcal{A}_{\Omega^- \Xi \pi}^{(P)} k_\mu u_\Omega^\mu \equiv G_F m_\pi^2 \frac{\alpha_{\Omega^- \Xi}^{(P)}}{\sqrt{2} f_\pi} \bar{u}_\Xi k_\mu u_\Omega^\mu ,$$

$$\begin{aligned} \Delta(\Xi^0 \pi^-) &\equiv \frac{\Gamma(\Omega^- \rightarrow \Xi^0 \pi^-) - \Gamma(\bar{\Omega}^- \rightarrow \bar{\Xi}^0 \pi^+)}{\Gamma(\Omega^- \rightarrow \Xi^0 \pi^-) + \Gamma(\bar{\Omega}^- \rightarrow \bar{\Xi}^0 \pi^+)} \\ &\approx \sqrt{2} \frac{\alpha_3^{(\Omega)}}{\alpha_1^{(\Omega)}} \sin(\delta_3 - \delta_1) \sin(\phi_3 - \phi_1) , \end{aligned}$$

P-wave dominance (predominantly parity conserving)

SM $\sim 2 \times 10^{-5}$ much larger than other **rate** asymmetries

Beyond SM could be 10-100(?) times larger

Tandean 04: $\Delta_{\Omega \rightarrow \Lambda \mathbb{K}} < 10^{-3}$ (but new constraints now)

Why is it larger in SM?

- The $\Delta I=3/2 / \Delta I=1/2$ amplitude ratio is about 7% from data (50% error?)
- The strong $\Xi\pi$ scattering phases estimated in lowest order χ PT $(\delta_1-\delta_3)\sim -14^\circ$ (error?)
- The weak phases are estimated in vacuum saturation to be $(\phi_1-\phi_3)\sim 0.001$ (tree operator suppressed, phase closer to phase of penguin \gg other hyperon decays)

from Tandean, Valencia **Phys.Lett.B451:382-387,1999,**

Phys.Lett.B452:395-401,1999